

# ***How to help keV right-handed neutrino dark matter?***

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***@5th International Workshop on Dark Matter, Dark Energy and Matter-Antimatter Asymmetry***

# Introduction

- **To consider new physics,**
  - **Neutrino masses**
  - **Dark matter candidate**
  - **Matter–antimatter asymmetry**

# Introduction

- **To consider new physics,**
  - **Neutrino masses**
  - **Dark matter candidate**
  - **Matter–antimatter asymmetry**
- + concept of minimality**



**the  $\nu$ MSM!**

# The $\nu$ MSM [Asaka, Blanchet, Shaposhnikov (2005)] [Asaka, Shaposhnikov (2005)]

- **The most economical extension : SM + 3RH $\nu$ s**

$$\mathcal{L}_N = i\bar{\nu}_{RI}\not{\partial}\nu_{RI} - F_{\alpha I}\bar{L}_{\alpha}H\nu_{RI} - \frac{M_N}{2}\overline{\nu_{RI}^c}\nu_{RI} + \text{h.c.}$$

## \*Key assumption

**Dirac masses :  $M_D \ll$  Majorana masses :  $M_N \lesssim \Lambda_{\text{EW}}$**

**\*Seesaw mechanism** [Minkowski (1977);Yanagida(1979);Gell-Mann,Ramond,Slansky (1979);  
Glashow (1980);Mohapatra,Senjanovic(1980)]

$$M_{\nu} = -M_D M_N^{-1} M_D^T \rightarrow 0.1 \text{ eV} \left( \frac{F}{10^{-7}} \right)^2 \left( \frac{1 \text{ GeV}}{M_N} \right)$$

**Tiny neutrino masses can still be explained  
by the seesaw mechanism with **small Yukawa coupling****

# The $\nu$ MSM

[Asaka, Blanchet, Shaposhnikov (2005)]

[Asaka, Shaposhnikov (2005)]

- Mass spectrum of  $RH\nu$ s

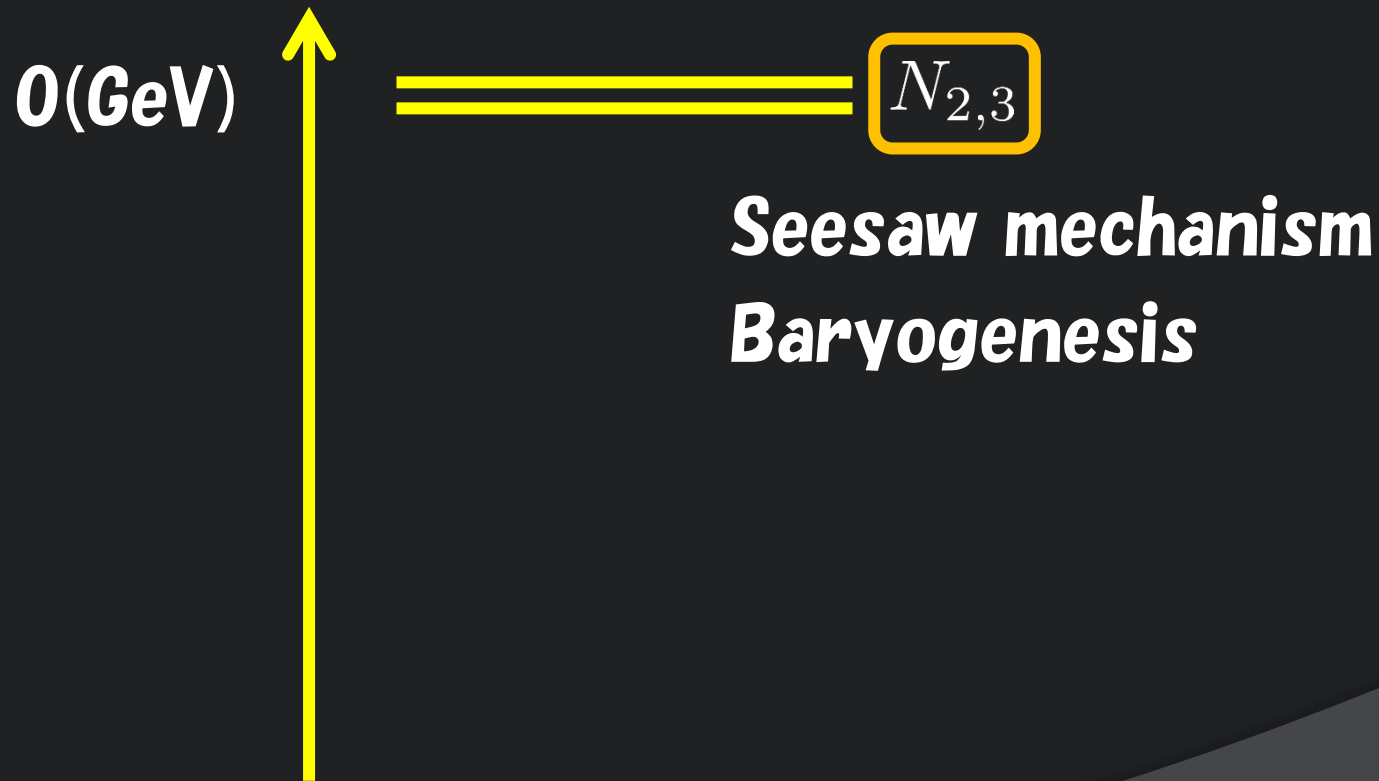


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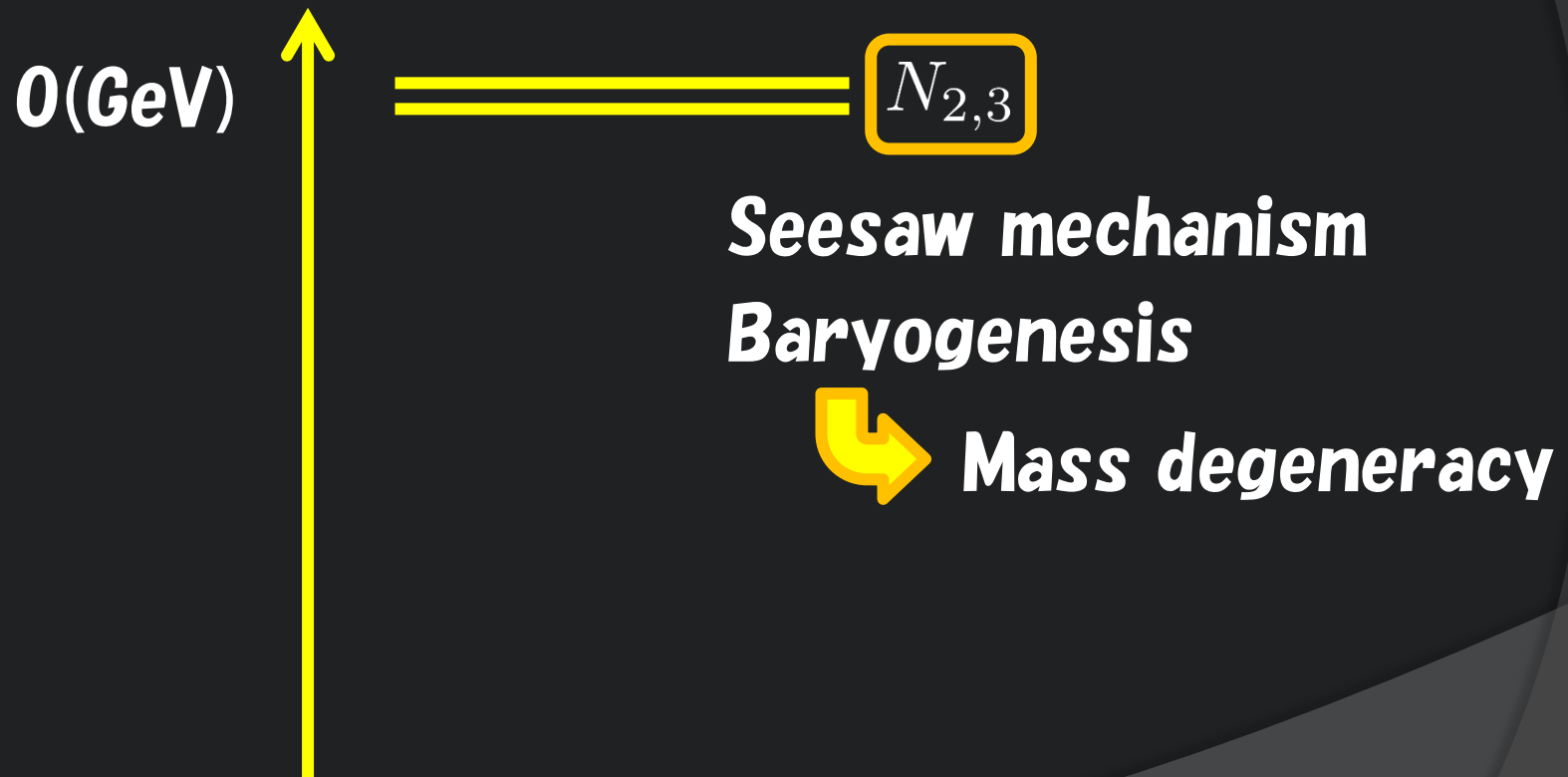


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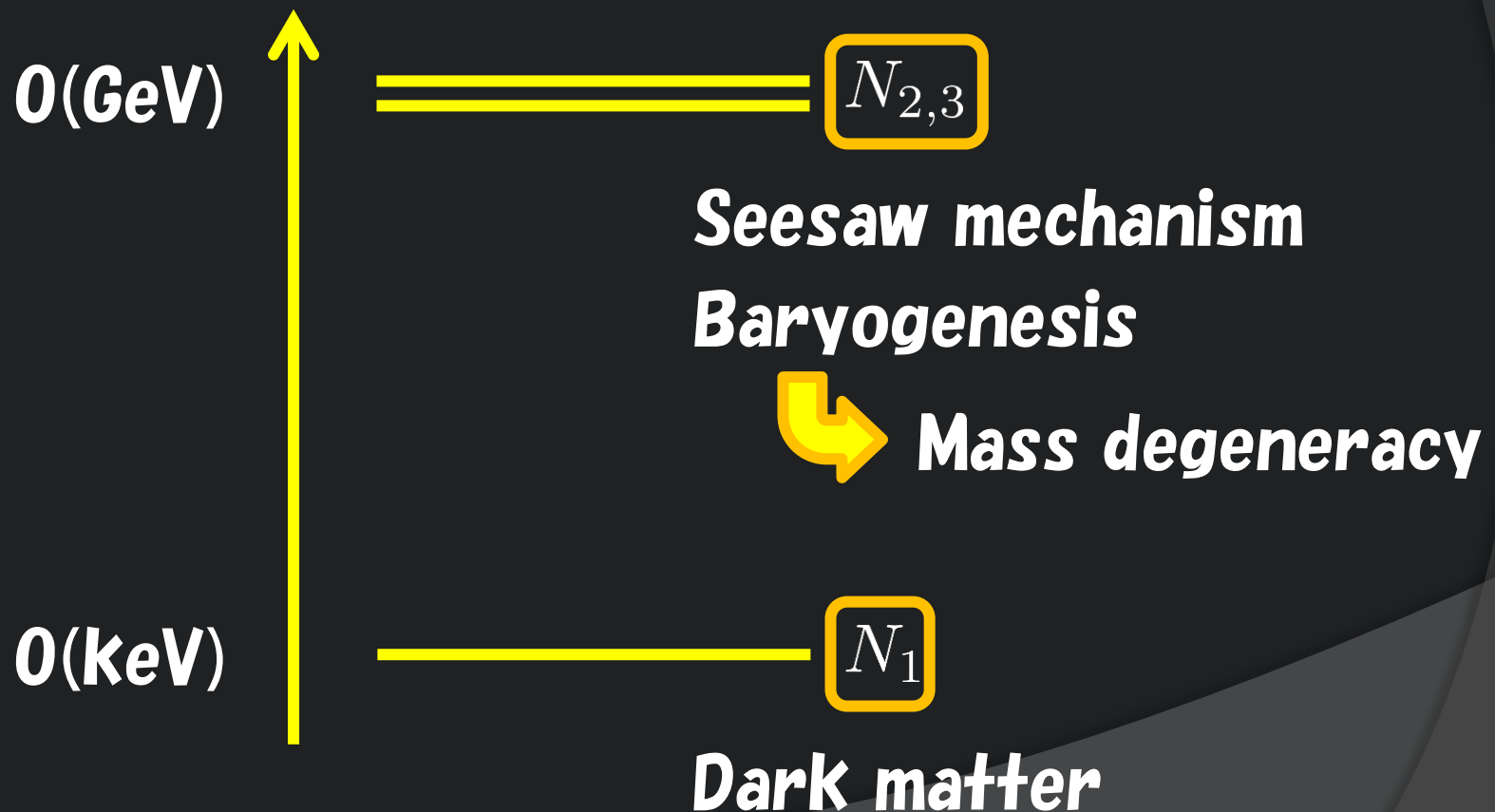


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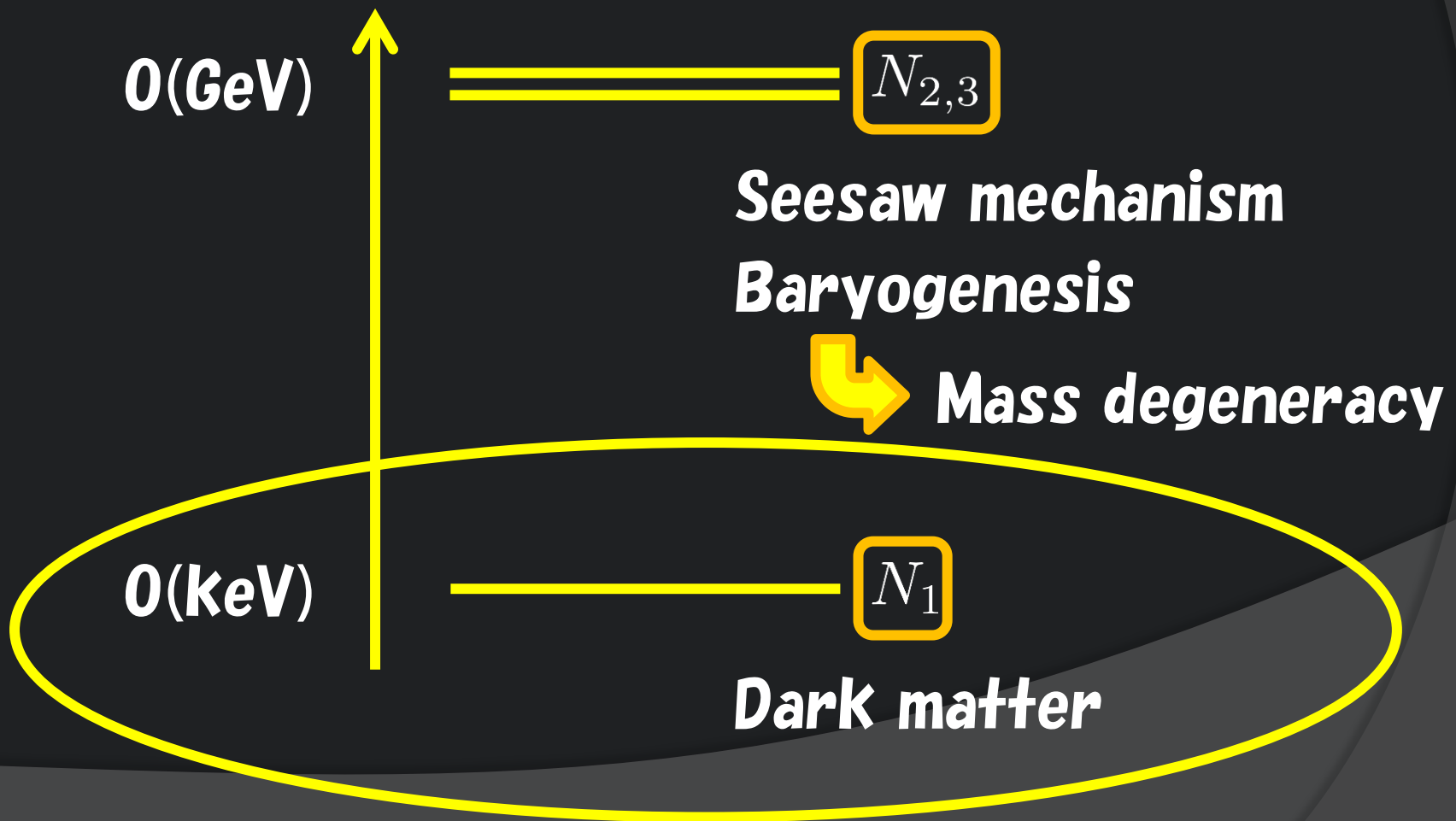


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# The $\nu$ MSM [Asaka, Blanchet, Shaposhnikov (2005)] [Asaka, Shaposhnikov (2005)]

## • *Physical states of neutrinos*

\* **active neutrinos** :  $\nu_i = U_{\text{MNS}}^\dagger \nu_{L\alpha} - U_{\text{MNS}}^\dagger \Theta \nu_{RI}^C$

\* **sterile neutrinos** :  $N_I^C = \nu_{RI}^C + \Theta^\dagger \nu_{L\alpha}$

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**right-handed neutrinos**

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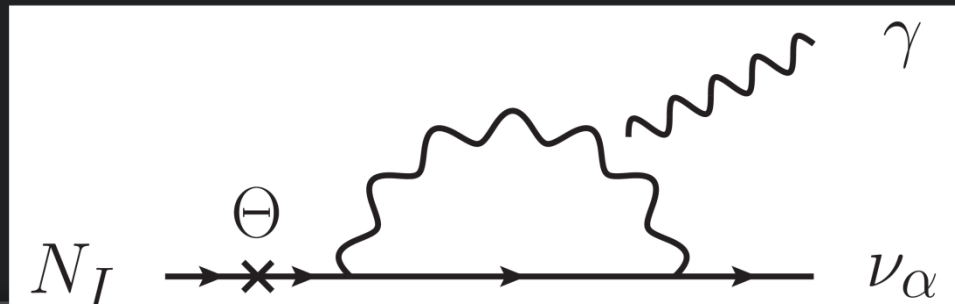
\* **sterile neutrinos** :  $N_I^C = \nu_{RI}^C + \Theta^\dagger \nu_{L\alpha}$



**right-handed neutrinos**

• **Important parameter** :  $\Theta \equiv M_D/M_N$

**Through this mixing,**



# The $\nu$ MSM

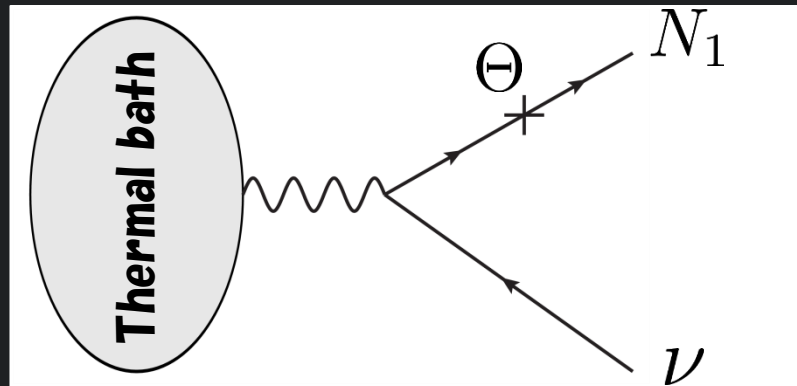
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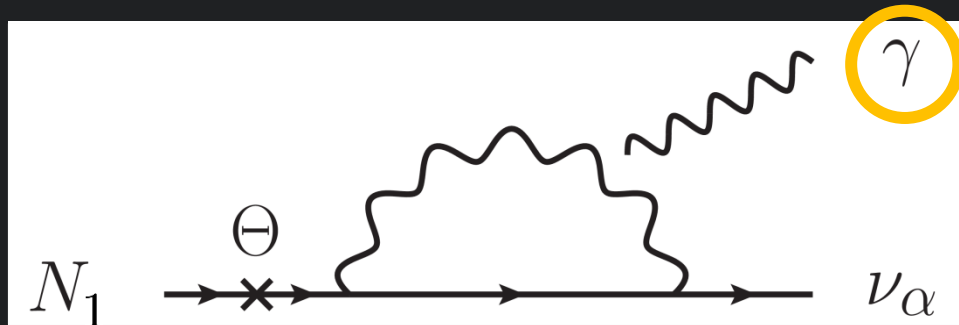
## • DM in the $\nu$ MSM

\*simple production mechanism (Dodelson–Widrow)

[Dodelson and Widrow (1993)]



– Mixing  – Production rate  



X-ray  
observatory

# The $\nu$ MSM

[Asaka, Blanchet, Shaposhnikov (2005)]

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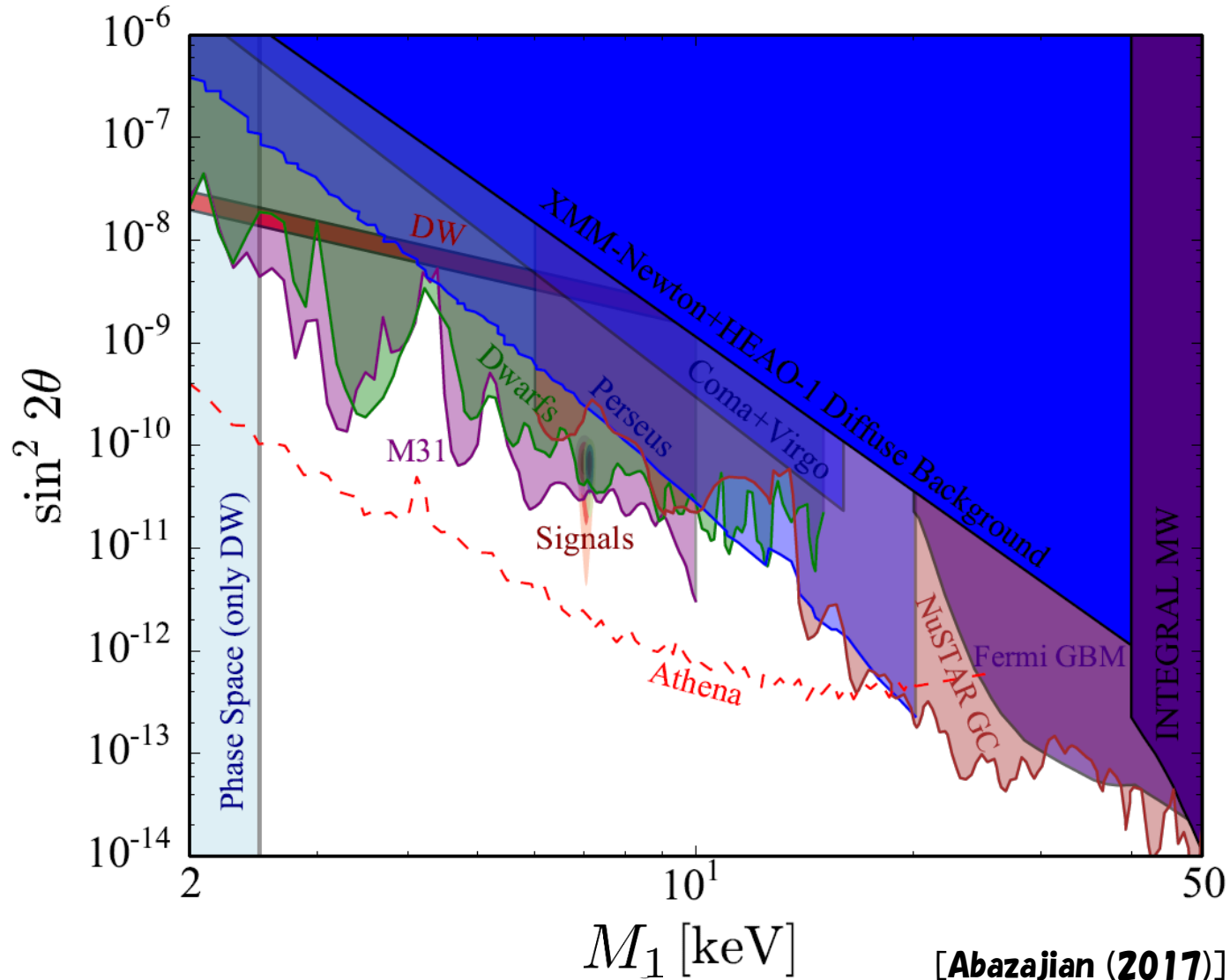
- *Constraints on DM*

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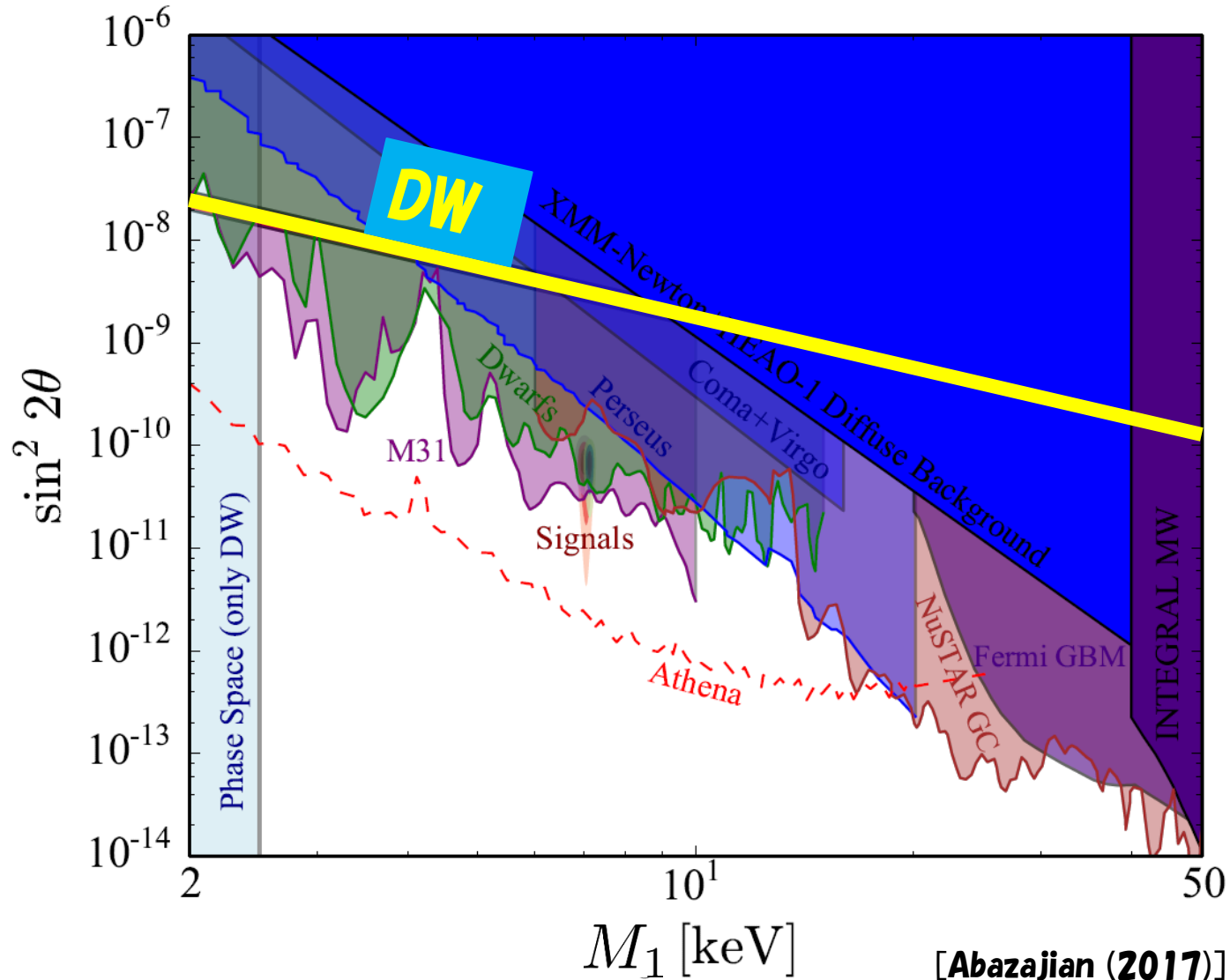


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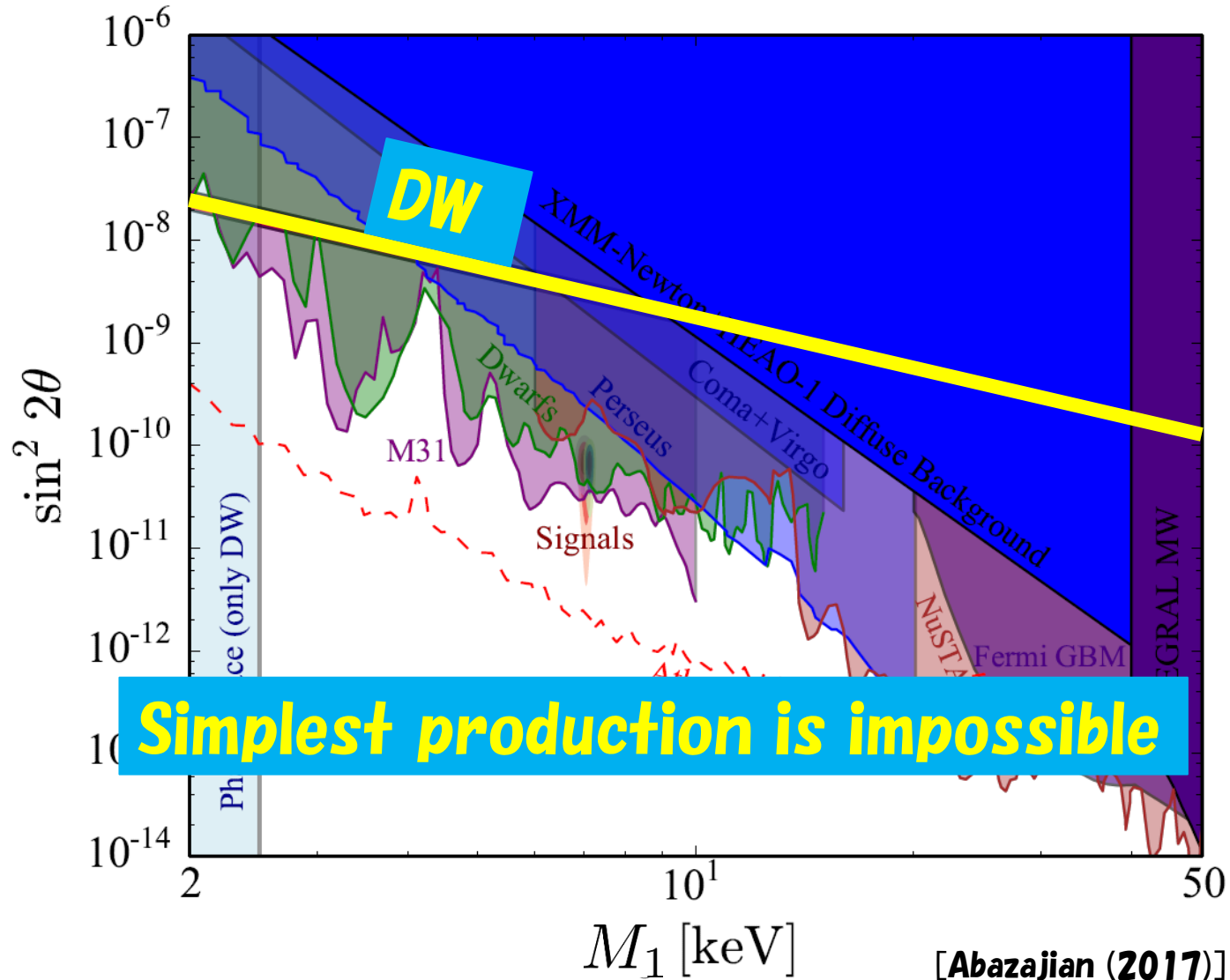


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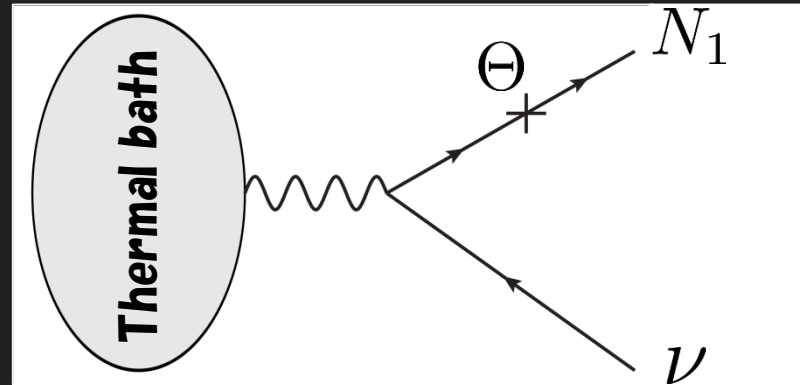
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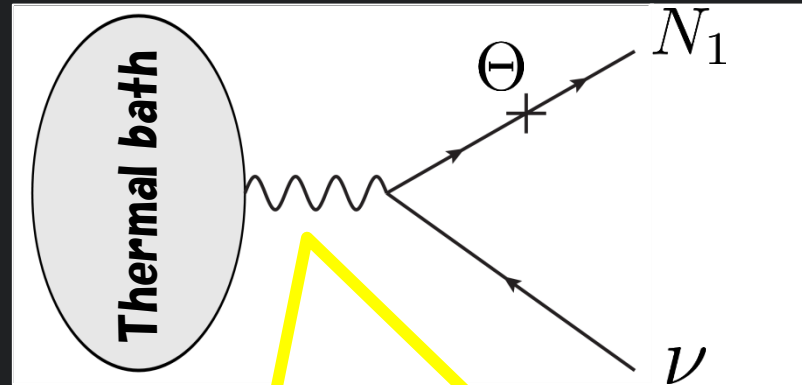
# keV DM production

- **Resonant production** [Shi and Fuller (1998)]



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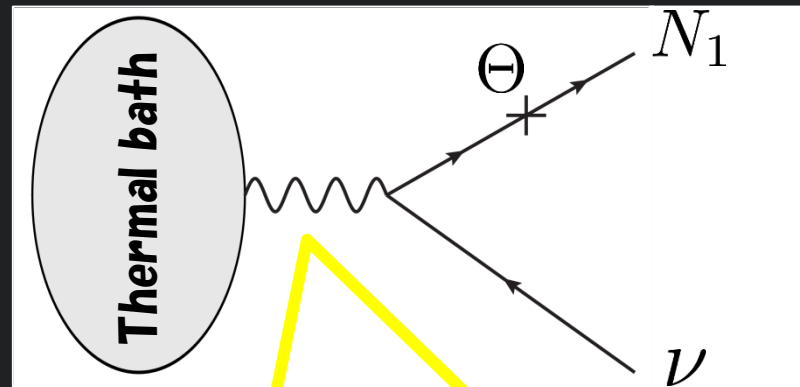
- **Resonant production** [Shi and Fuller (1998)]



**Large # of lepton asymmetry**

# keV DM production

- **Resonant production** [Shi and Fuller (1998)]



**Large # of lepton asymmetry**

**\*typical required amount of Lepton asymmetry**

$$\eta_L \equiv n_L/n_\gamma \gtrsim 10^6 \eta_B$$

**How to produce?**

# **keV DM production**

- **Late time resonant leptogenesis**

**\*same idea as one leptogenesis for baryogenesis**

[Pilaftsis and Underwood (2003, 2005)]

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- small Yukawa coupling**

- mass degeneracy of  $RH_\nu$ s**

# keV DM production

- Late time resonant leptogenesis

- \* same idea as one leptogenesis for baryogenesis

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resonant lepton asymmetry production  
after EWSB is possible!

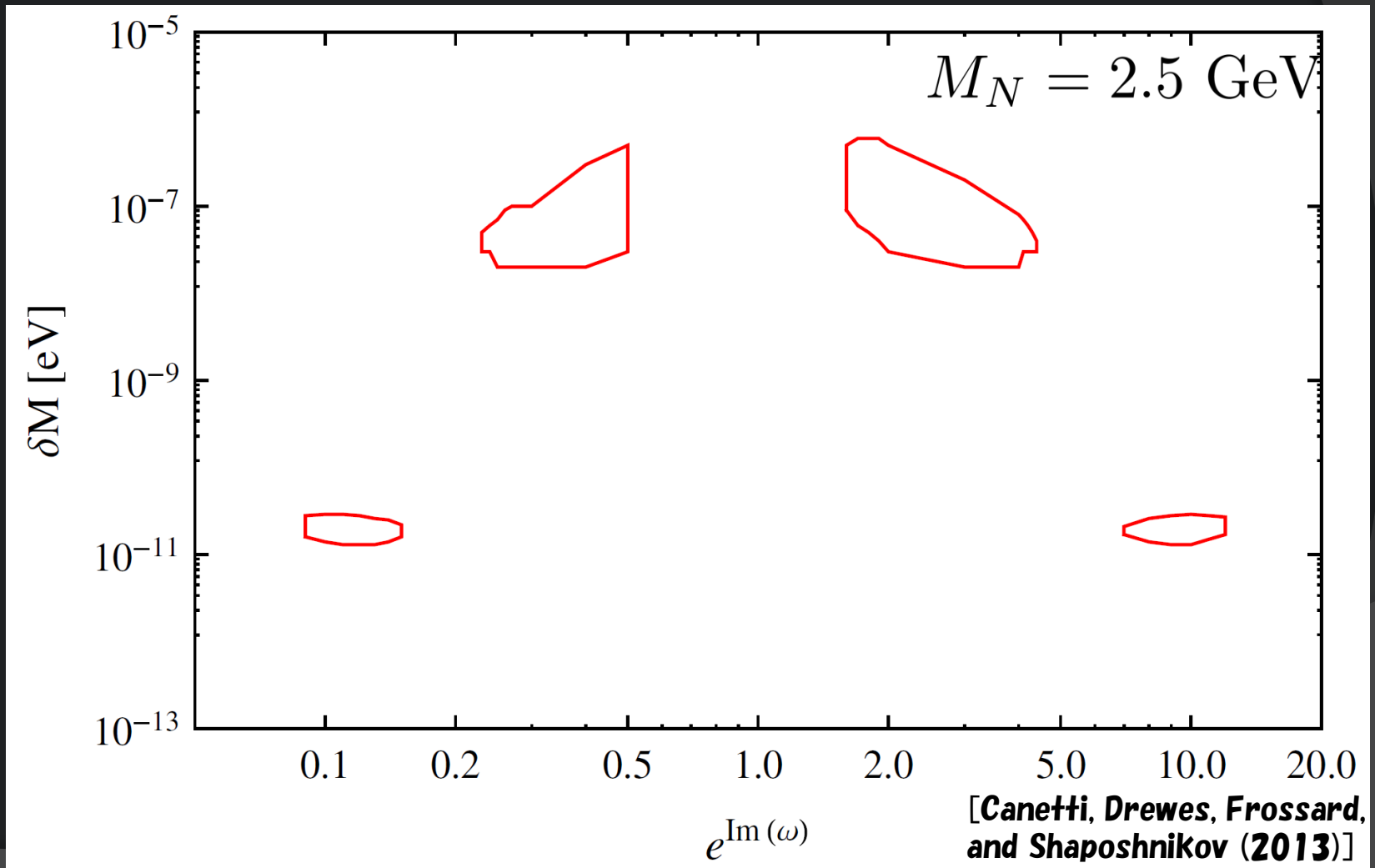
# ***keV DM production***

- ***Late time lepton asymmetry production***



# keV DM production

- Late time lepton asymmetry production

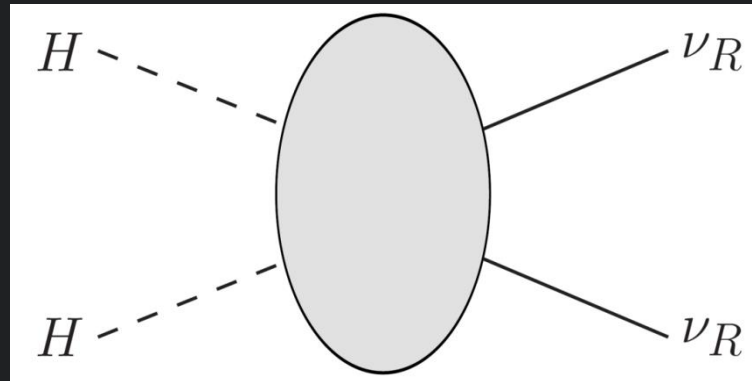


# keV DM production

- **Higher dimensional production @high scale**

[Asaka, Eijima, Hl, Minogawa, Yoshii (2017)]

$$\mathcal{L}_{\text{HD}} = \frac{A_{IJ}}{\Lambda} H^\dagger H \overline{\nu_{RI}^c} \nu_{RJ} + \text{h.c.}$$



## Produced amount of RH $\nu$ s

$$\begin{aligned} [\rho_N^I]_{IJ} &= \# \times \frac{M_P T_R}{\Lambda^2} [A^\dagger A]_{IJ} \\ &\xrightarrow{\Lambda \rightarrow M_P} \# \times \frac{T_R}{M_P} [A^\dagger A]_{IJ} \end{aligned}$$

# Higher dimensional operator effect

[Asaka, Eijima, Hl, Minogawa, Yoshii (2017)]

## • Effect on (constraint from) DM production

If 
$$\sum_{\alpha} |\Theta_{\alpha 1}|^2 \ll 8 \times 10^{-8} \left( \frac{M_1}{1 \text{ keV}} \right)^{-2}$$

 **DW mechanism does not work at all**

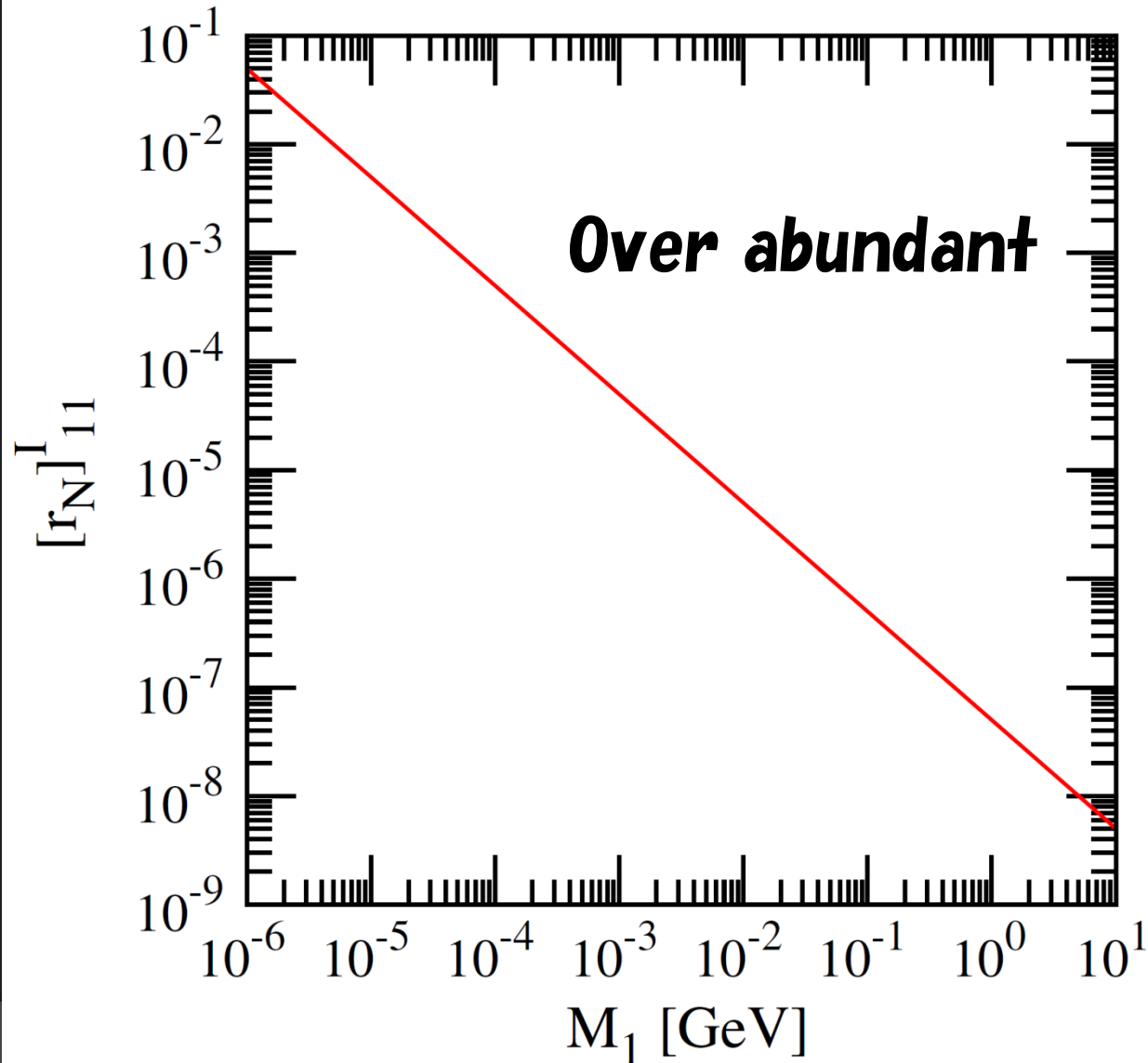
 **Production via HD operators becomes more important!!**

# Higher dimensional operator effect

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- Effect on (constraint from) DM production

If



at all

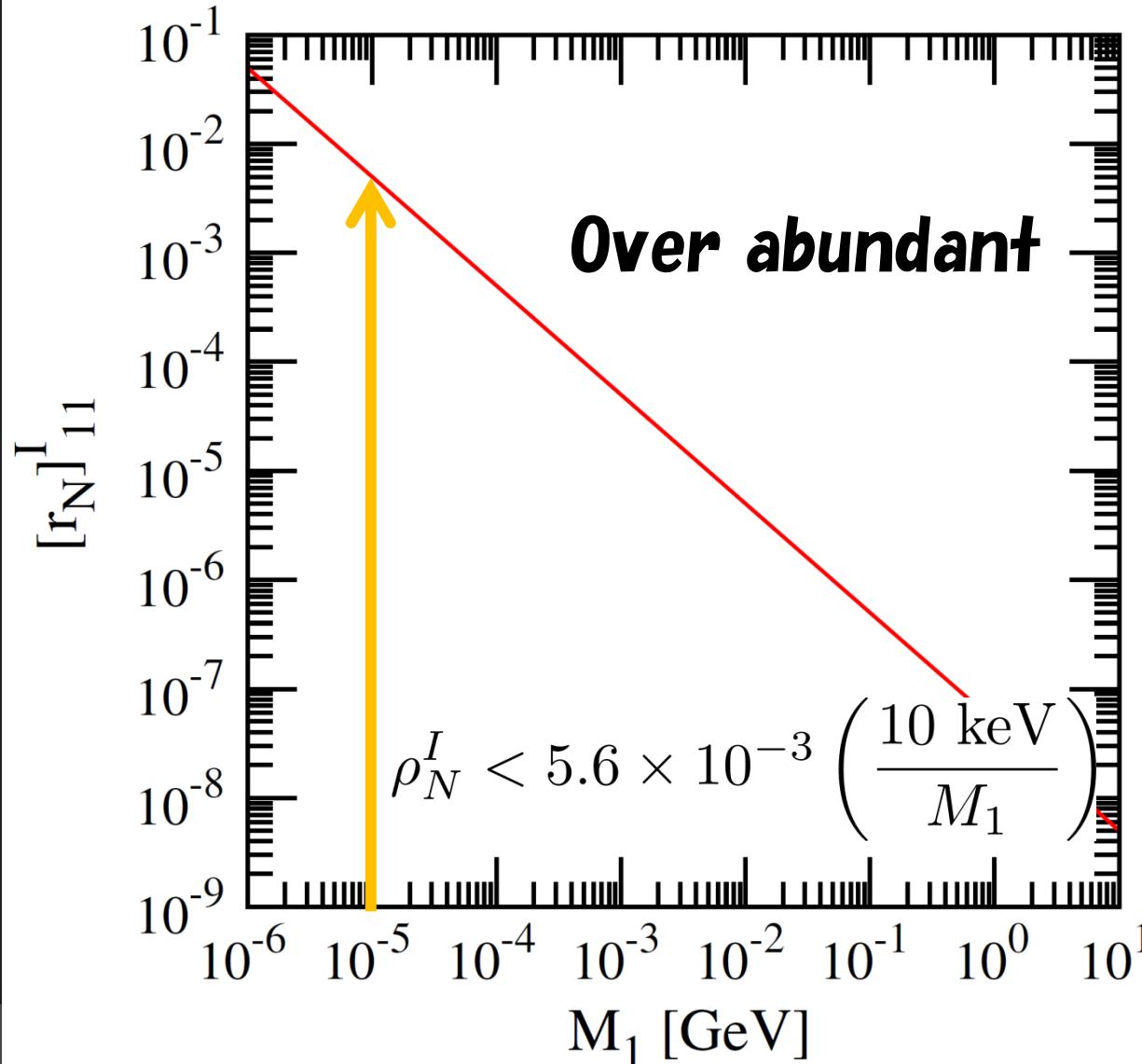
s

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# ***Higher dimensional operator effect***

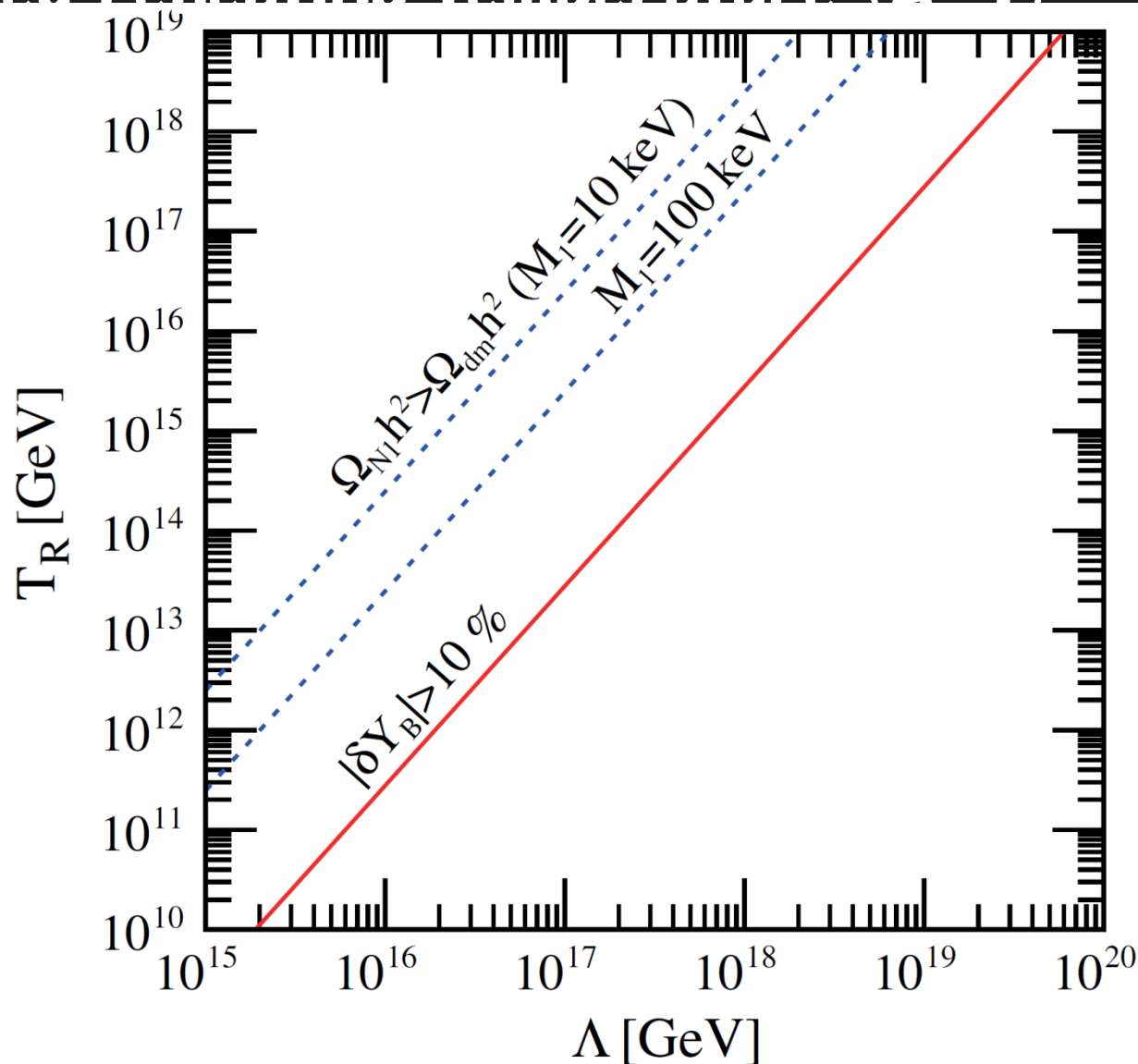
[Asaka, Eijima, Hl, Minogawa, Yoshii (2017)]

- ***Predicted reheating temperature ( $A=1$ )***

# Higher dimensional operator effect

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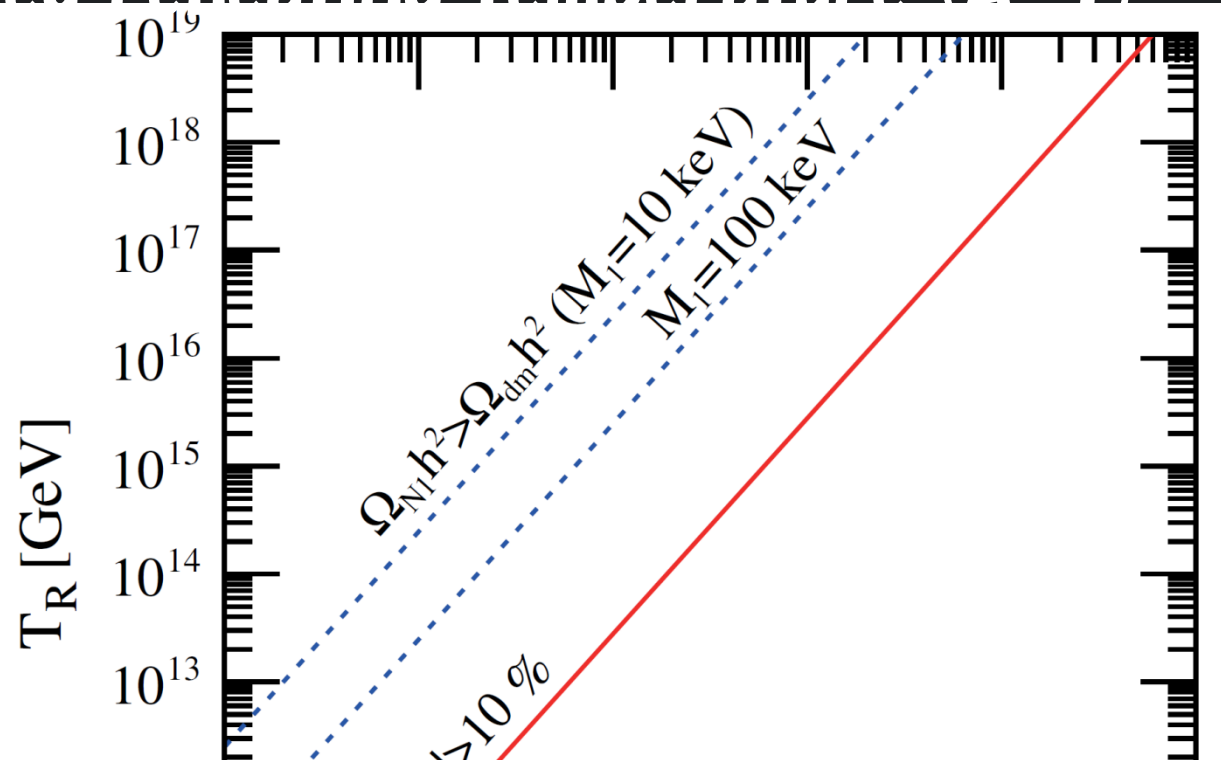
## • Predicted reheating temperature ( $\Lambda = 1$ )



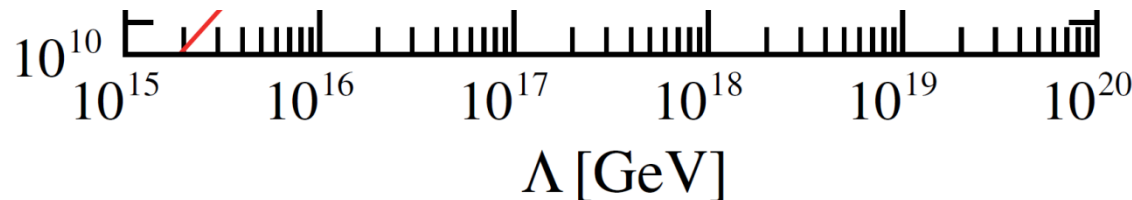
# Higher dimensional operator effect

[Asaka, Ejima, Hl, Minogawa, Yoshii (2017)]

## • Predicted reheating temperature ( $\Lambda = 1$ )



**Once DM mass and reheating temperature are given coupling strength of HD operator can be predicted!**

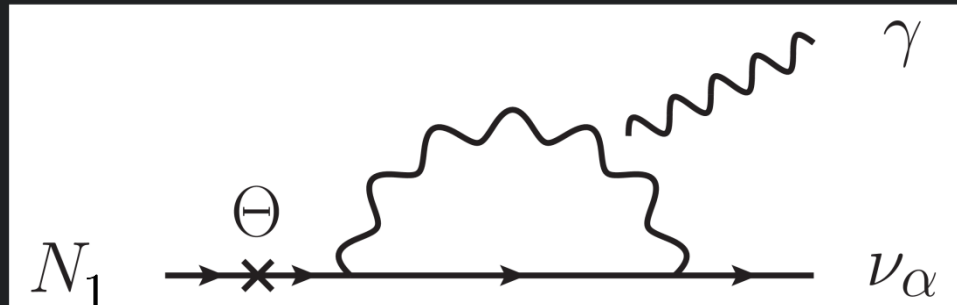




# Higher dimensional operator effect

[Asaka, Eijima, Hl, Minogawa, Yoshii (2017)]

- **When**  $T_R = 10^{13}$  GeV,  $A = \mathcal{O}(1)$ ,  $\Lambda = M_P$   
 $M_{N_1} \simeq 5$  GeV

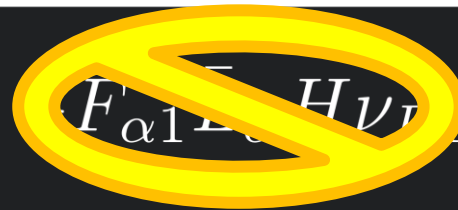
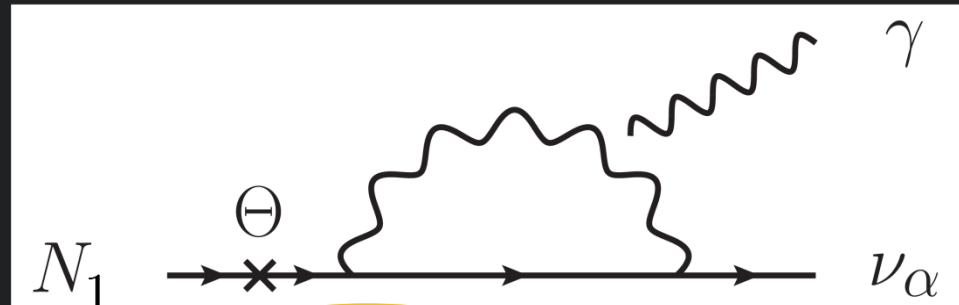


$$-F_{\alpha 1} \bar{L}_{\alpha} H \nu_{R1}$$

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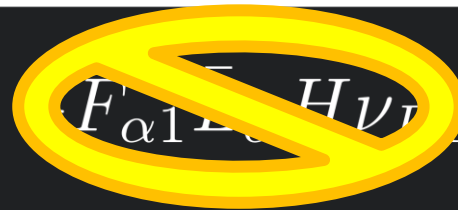
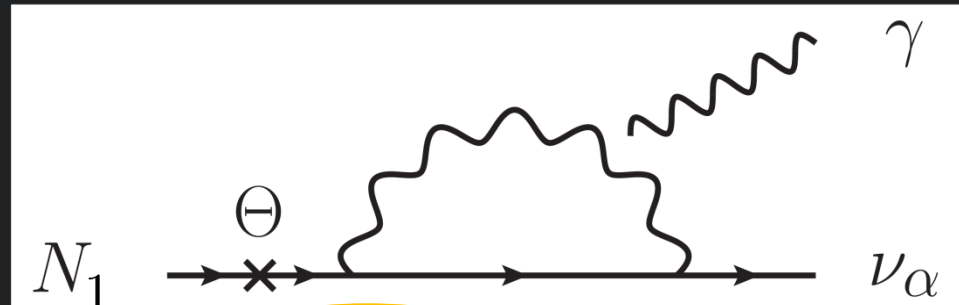


$\mathbf{Z}_2$  need to be imposed

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$Z_2$  need to be imposed

$\frac{A_{11}}{\Lambda} H^\dagger H \overline{\nu_{R1}^c} \nu_{R1}$  is invariant under  $Z_2$

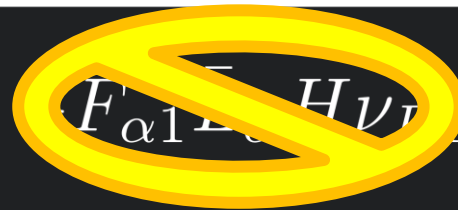
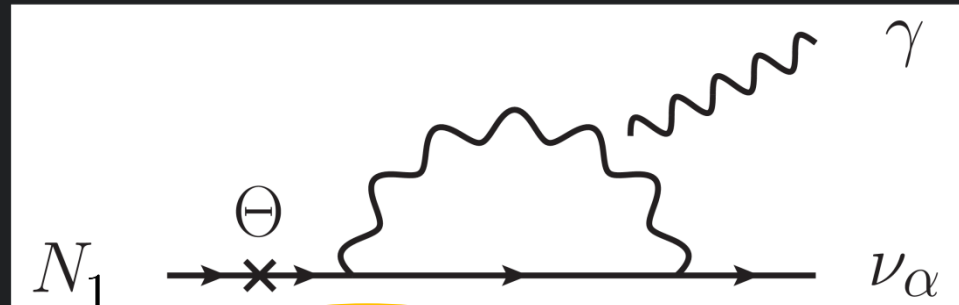


**This DM production becomes important!**

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$F_{\alpha 1} \bar{L}_\alpha H \nu_{R1}$   $\mathbf{Z}_2$  need to be imposed

$\frac{A_{11}}{\Lambda} H^\dagger H \overline{\nu_{R1}^c} \nu_{R1}$  is invariant under  $\mathbf{Z}_2$

➡ This DM production becomes important!

Test might become difficult...

# Summary

- **keV right-handed neutrino is still one of the possibilities of DM candidate**

- \*Simplest production does not work**

- Resonant production**

- Highly tuned parameters are required**

- Higher dimensional operator production**

- What is the underlying theory?**

- Welcome to discuss about alternative possibilities!**

***Thank you for  
your attention!***

